

# Computing plates and some related topics

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## Outline of the talk

- Motivation
- Some formulations of the problem
- Analytic properties of a mixed formulation
- Reproducing kernels and the biharmonic operator
- Positivity issues in the plate bending problems
- Some properties of the systems of linear algebraic equations obtained via discretization

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# 1 Definitions and notation

$\Omega$  - region in  $\mathcal{R}^2$

$\partial\Omega$  - boundary of region  $\Omega$

$H^m(\Omega)$ ,  $1 \leq m \leq 2$  - solution space

$\|v\|_{m,\Omega} = \left( \sum_{|\alpha| \leq m} \int_{\Omega} |\partial^\alpha v|^2 \right)^{1/2}$  - norm in the space  $H^m(\Omega)$

$H_0^m(\Omega)$  - zero-trace solution space

$L^{1/2}(\partial\Omega)$  - trace space

- $\mathcal{M}$  - complementary subspace of the Hilbert space  $H^1(\Omega)$  with respect to  $H_0^1(\Omega)$   
thus satisfies relation  $H^1(\Omega) = H_0^1(\Omega) \oplus \mathcal{M}$
- $(\cdot, \cdot)_{\mathcal{M}}$  - inner product on  $\mathcal{M}$  inducing on  $\mathcal{M}$  norm equivalent with the norm  
heredited from  $H^1(\Omega)$

## 2 Biharmonic equation

$$\Delta^2 u = f \quad \text{in } \Omega \quad \text{and} \quad u = 0 = \frac{\partial u}{\partial \nu} \quad \text{on } \partial\Omega$$

### 2.1 Formulation of the Problem

Find  $u \in H_0^2(\Omega)$  such that relations

$$(2.1) \quad J(u) = \min \{ J(v) : v \in H_0^2(\Omega) \},$$

hold where

$$(2.2) \quad J(v) = \frac{1}{2} \int_{\Omega} |\Delta v|^2 - \int_{\Omega} f v, \quad v \in H_0^2(\Omega).$$

In place of (2.2) one can minimize functional

$$(2.3) \quad \mathcal{J}(v, \psi) = \frac{1}{2} \int_{\Omega} |\psi|^2 - \int_{\Omega} f v$$



under the assumption that  $v \in H_0^2$  and  $\psi \in L^2(\Omega)$  and simultaneously

$$\Delta v = \psi.$$

The subspaces corresponding to the variational principle introduced are characterized as follows:

$$\mathcal{V} = \{(v, \psi) \in H_0^1(\Omega) \times L^2(\Omega) : \forall \mu \in H_0^1(\Omega) \text{ satisfying } \beta((v, \mu), \mu) = 0\},$$

where

$$(2.4) \quad \beta((v, \psi), \mu) = \int_{\Omega} \text{grad} v \text{ grad} \mu - \int_{\Omega} \psi \mu$$

A relationship between the natural variational principle (2.1)-(2.2) and that one given by (2.3) reads:

**2.2** *Assume  $u$  is a solution of Problem 1. Then also*

$$\mathcal{J}(u, -\Delta u) = \min \{\mathcal{J}(v, \phi) : (v, \phi) \in \mathcal{V}\}.$$

This fact guarantees a possibility to factorize Problem 1 to a subsequent solving Poisson problems, hence problems of order 2. It is easy to see that function  $\phi$  can be interpreted as an approximation or representative of  $-\Delta u$ .

### 3 An Algorithm

**Input data:**  $f \in L^2(\Omega)$ ,  $\lambda^0 \in \mathcal{M}$

1<sup>(0)</sup> Set subsequently  $k = 0, 1, \dots$

2<sup>0</sup> Find  $\phi^k \in H_0^1(\Omega)$  such that in the classical formulation

$$\Delta \phi^k = f \text{ in } \Omega$$

$$\phi^k = \lambda^k \text{ on } \Gamma$$

in the variational formulation

$$\int_{\Omega} \text{grad} \phi^k \text{grad} \mu d\Omega = \int_{\Omega} f \mu d\Omega \quad \forall \mu \in H_0^1(\Omega)$$

$$\phi^k - \lambda^k \in H_0^1(\Omega)$$

3<sup>0</sup> Find  $u^k \in H_0^1(\Omega)$  such that  
in the classical formulation

$$\Delta u^k = \phi^k \text{ in } \Omega$$

$$u^k = 0 \text{ on } \partial\Omega$$

in the variational formulation

$$\int_{\Omega} \text{grad} u^k \text{grad} \mu d\Omega = \int_{\Omega} \phi^k \mu d\Omega \quad \forall \mu \in H_0^1(\Omega)$$

$$u^k \in H_0^1(\Omega)$$

4<sup>0</sup> Find  $\lambda^{k+1} \in \mathcal{M}$  such that  
in the classical formulation

$$\lambda^{k+1} = \lambda^k + \rho [\Delta u^k - \phi^k] \text{ on } \partial\Omega$$

in the variational formulation

$$(\lambda^{k+1} - \lambda^k, \mu)_{\mathcal{M}} = \rho \left[ \int_{\Omega} \text{grad} u^k \text{grad} \mu d\Omega - \int_{\Omega} \phi^k \mu d\Omega \right]$$

$$\forall \mu \in \mathcal{M}$$

Convergence of the Algorithm is characterized using the following relations.

3.1 Assuming parameter  $\rho$  is chosen in the interval  $(0, 2c^2\sigma^2)$ , where

$$\sigma = \inf \left\{ \frac{\|\Delta v\|_{L^2(\Omega)}}{\|\frac{\partial v}{\partial \nu}\|_{L^2(\Omega)}} : v \in H^2(\Omega) \cap H_0^1(\Omega) \right\}$$

and constant  $c > 0$  satisfies relations

$$c\|\mu\|_{L^2(\partial\Omega)} \leq (\mu, \mu)_{\mathcal{M}}^{1/2}, \quad \forall \mu \in H^1(\Omega),$$

the quantities appearing in Algorithm 1 fulfil relations

$$\lim_{k \rightarrow \infty} \|u^k - u\|_{1,\Omega} = 0,$$

a

$$\lim_{k \rightarrow \infty} \|\phi^k + \Delta u\|_{L^2(\Omega)} = 0.$$

## 4 Discrete formulation

Space of approximate solutions:

$$V_h \subset H^1(\Omega)$$

Space of approximate solutions with zero traces:

$$V_{0h} = \{v_h \in V_h : v_h = 0 \text{ na } \partial\Omega\}$$

Space of approximate solutions with constraints:

$$\mathcal{V}_h = \{(v_h, \psi_h) \in V_{0h} \times V_h; \forall \mu_h \in V_h, \beta((v_h, \psi_h), \mu_h) = 0\}$$

Space of approximate solutions with constraints and zero traces:

$$\mathcal{W}_h = \{(v_h, \psi_h) \in V_{0h} \times V_h; \forall \mu_h \in V_{0h}, \beta((v_h, \psi_h), \mu_h) = 0\}$$

Complementary space of  $V_{0h}$  in  $V_h$ , thus,

$$V_h = V_{0h} \oplus \mathcal{M}_h$$



#### 4.1 Algorithm $1_h$

Input data:  $f_h \in V_h$ ,  $\lambda_h^0 \in \mathcal{M}$

Set subsequently  $k = 0, 1, \dots$

1<sup>0</sup> Find  $\phi_h^k$  satisfying relations  $\phi_h^k - \lambda_h^k \in V_{0h}$  and

$$\forall v_h \in V_{0h}, \quad \int_{\Omega} \text{grad } \phi_h^k \text{ grad } v_h = \int_{\Omega} f v_h.$$

2<sup>0</sup> Find  $u_h^k \in V_{0h}$  such that by

$$\forall v_h \in V_{0h}, \quad \int_{\Omega} \text{grad } u_h^k \text{ grad } v_h = \int_{\Omega} \phi_h^k v_h.$$

3<sup>0</sup> Find  $\lambda_h^{k+1} \in \mathcal{M}_h$  such that

$$\forall \mu_h \in \mathcal{M}_h, \quad (\lambda_h^{k+1} - \lambda_h^k, \mu_h)_{\mathcal{M}_h} = \rho \beta ((u_h^k, \phi_h^k), \mu_h) = 0.$$

Let  $A_h : V_h \rightarrow V_{0h}$  be given by

$$A_h \psi_h = v_h \Leftrightarrow \forall \mu_h \in V_{0h}, \int_{\Omega} \text{grad } v_h \text{ grad } \mu_h = \int_{\Omega} \psi_h \mu_h$$

Similarly,  $B_h : V_h \rightarrow \mathcal{M}_h$  utilizing the system of equations:  $B_h$

$$\forall \mu_h \in \mathcal{M}_h : (B_h \psi_h, \psi_h)_{\mathcal{M}_h} = \beta ((A_h \psi_h, \psi_h), \mu_h) = 0.$$

Hence,  $B_h \psi_h$  is a restriction of  $A_h \psi_h$  onto  $\partial\Omega$ . (We have that  $A_h \sim \Delta^{-1}$  with boundary conditions implied from  $A_h \psi_h$ ).

A result similar to that describing convergence for the continuous case is contained in the following

**4.2** Assuming parameter  $\rho$  is chosen from the interval  $(0, 2\sigma_h^2)$ , where

$$\sigma_h = \frac{1}{\|B\|},$$

and  $\|B\|$  denotes the operator  $L^2$ -norm, relations

$$\lim_{k \rightarrow \infty} \|u_h^k - u_h\|_{V_{0h}} = 0, \quad \lim_{k \rightarrow \infty} \|\phi_h^k - \phi_h\|_{V_h} = 0$$

hold.

## 5 Finite element spaces

Triangulation  $\mathcal{T}_h$  of the region  $\Omega$  consists of elements  $K$  such that  $\bar{\Omega} = \bigcup_{K \in \mathcal{T}_h} K$  and each element  $K$  satisfies for all triangulations  $0 < h \leq h_0$  standard geometric conditions in the sense of [5]. It is assumed also that each element  $K$  is an affine picture of a reference element  $\hat{K} : K = F_K(\hat{K})$ .

Space

$$V_h = \{v_h \in \mathcal{C}(\bar{\Omega}) : \forall K \in \mathcal{T}_h, v_h|_K \in P_K\}$$

is generated utilizing the triangulation  $\mathcal{T}_h$  where

$$P_K = \left\{ v : K \rightarrow \mathcal{R}; v = \hat{v} \circ F_K^{-1}, \forall \hat{v} \in \hat{P} \right\}$$

assuming  $\hat{P}$  a finite dimensional space of functions satisfying

$$P_1 \subset \hat{P}$$

where  $P_1$  denotes the set of all polynomials of order  $\leq 1$  in two variables.

**5.1** As regular class of triangulations one understands a system  $\{\mathcal{T}_h\}$  satisfying with suitable constants independent of  $h$ ,  $\alpha$  and  $\tau$  such that

$$\max \left\{ \frac{h(K)}{\delta(K)} \right\} \leq \alpha,$$

$$\tau \max \{h(K) : K \in \mathcal{T}_h\} \leq \min \{h(K) : K \in \mathcal{T}_h\}, \quad h = \max \{h(K) : K \in \mathcal{T}_h\}$$

where  $h(K) = \text{diameter } K$ ,  $\delta(K) = \sup\{\text{diameter of the discs inscribed into } K\}$ .

The above theory has been developed without any requirements concerning the spaces  $V_h$ ,  $V_{0h}$ ,  $\mathcal{M}_h$ . To determining the value of parameter  $\rho$  we need some specifications concerning these spaces however. As spaces  $\mathcal{M}_h$  we choose in agreement with the choice of the spaces  $V_h$  a  $V_{0h}$  the following specific ones: Spaces  $\mathcal{M}_h$  are subspaces of  $V_h$  consisting of functions whose values vanish at all inner nodes of  $\Omega$ . We then have the following result.

**5.2 Theorem** Assume that the inner product  $(\cdot, \cdot)_{\mathcal{M}_h}$  is the  $L^2$ -inner product on  $\partial\Omega$ . Further, let  $V_h$  a  $V_{0h}$  a  $\mathcal{M}_h$  be chosen such described above.

Then

$$\lim_{k \rightarrow \infty} \sigma_h = \sigma,$$

where  $\sigma$  is the quantity introduced in the part of this paper devoted to the continuous problem.

## 6 Reproducing kernels and the biharmonic operator

Consider in a plane domain  $\mathcal{D}$  the class  $\mathcal{B}$  of all regular harmonic functions (in general complex valued) with a finite norm given by

$$\|h\|^2 = \int \int_{\mathcal{D}} |h(x, y)|^2 dx dy.$$

Class  $\mathcal{B}$  possesses a reproducing kernel which will be denoted by  $H(z, z_1)$ , i.e.

$$f(z) = \int \int_{\mathcal{D}} H(z, z_1) f(z_1) dz_1, \quad \forall f \in \mathcal{B}, z = x + iy \in \mathcal{D}.$$

Assume  $u$  is a complex-valued function. Then  $\hat{u}$  denotes its complex conjugate.

**6.1 Theorem [1]** *In order that the reproducing kernel  $K(x, y)$  of the proper functional Hilbert space  $\mathcal{E}$  be nonnegative it is necessary and sufficient that  $\mathcal{E}$  have two properties*

(a) *If  $u \in \mathcal{E}$ , then  $\hat{u} \in \mathcal{E}$  and  $\|\hat{u}\| = \|u\|$ .*

(b) *For each real-valued  $u \in \mathcal{E}$  there exists  $\tilde{u} \in \mathcal{E}$  such that*

$$\tilde{u}(x) \geq |u(x)| \text{ for all } x \text{ and } \|\tilde{u}\| = \|u\|.$$



## 7 Positivity issues in plate bending problems

Assume  $\mathcal{D} \subset \mathcal{R}^2$  is a region.

Denote by symbol  $A$  an operator defined by

$$\begin{cases} Au \equiv \Delta^2 u & \text{for } (x, y) \in \mathcal{D} \\ u(x, y) = 0 = \frac{\partial}{\partial \nu} u(x, y) & \text{for } (x, y) \in \partial \mathcal{D} \end{cases}$$

where

$$(\Delta u)(x, y) \equiv \left[ \left( \frac{\partial}{\partial x} \right)^2 + \left( \frac{\partial}{\partial y} \right)^2 \right] u(x, y) \quad \text{for } (x, y) \in \mathcal{D}.$$

**Question (Hadamard)** Does the pointwise relation

$$Au(x, y) \geq 0 \quad \text{for all } (x, y) \in \mathcal{D}$$

imply that

$$u(x, y) \geq 0 \quad \text{for all } (x, y) \in \mathcal{D}?$$

**Answer Not always!**

According to Duffin [7] the answer to the above question is **no** if  $\mathcal{D}$  is suitably chosen rectangle.

The answer to the Hadamard question is **affirmative** if

- $\mathcal{D}$  is a disc

Moreover, Duffin conjectures [7] that the answer to the Hadamard question is affirmative if  $\mathcal{D}$  is a square.

## 8 Concluding remarks

- The original problem of the fourth order is transferred to iterating on two Poisson equations in the domain  $\Omega$  and finding solution to one functional equation on the boundary  $\partial\Omega$ .
- The structure of the Poisson problems to solve is suitable for application of the most efficient (multigrid) methods: Actually, the corresponding matrices are symmetric  $M$ -matrices
- The operator governing the functional equation on the boundary  $\partial\Omega$  reads

$$\Delta^{-1} P_{\mathcal{M}, H_0^1(\Omega)} \Delta^{-1} \Delta$$

where  $P_{\mathcal{M}, H_0^1(\Omega)}$  denotes the projection of  $H^1(\Omega)$  onto  $\mathcal{M}$  along  $H_0^1(\Omega)$  and is thus symmetrizable and positive semidefinite. A simple Richardson iteration recommended in [5] to compute approximate boundary data for the auxiliary Poisson prob-

lem can, if needed, be replaced by some efficient Krylov like method, e.g. *conjugate gradient method*.

- An extra additional gain of the approach described is continuity of the quantity  $\Delta u$ , where  $u$  is the true solution of our biharmonic problem. This property of  $\Delta u$  may not be guaranteed even if one applies some conformal FEM.
- In order to use the methods described in the lecture to computing plate problems standard software products utilizing  $C^0$  elements apply .